

Coalescence of phospholipid vesicles mediated by β 2GPI – experiment and modelling

J. Urbanija¹, B. Rozman³, A. Iglič², T. Mareš⁴, M. Daniel⁴, Veronika Kralj-Iglič¹

¹Laboratory of Clinical Biophysics, Faculty of Medicine, University of Ljubljana, Lipičeva 2, Ljubljana, Slovenia;

²Laboratory of Physics, Faculty of Electrical Engineering, University of Ljubljana, Tržaška 25, Ljubljana, Slovenia;

³Department of Rheumatology, University Medical Centre, Vodnikova 62, Ljubljana, Slovenia

⁴Laboratory of Biomechanics, Faculty of Mechanical Engineering, Czech Technical University in Prague, Prague, Czech Republic

Abstract— Collective interactions between the giant phospholipid vesicles made of POPC, cardiolipin and cholesterol after the addition of β 2GPI may cause the coalescence of membrane buds to the mother cell. Using the discrete elastic model of the vesicle membrane mechanics it was shown that the coalescence of the buds depends on the adhesion strength and rigidity of the biomembrane.

Keywords— Beta-2-glycoprotein-I; Apolipoprotein, elastic model, discretization.

I. INTRODUCTION

The serum protein β 2GPI is considered to have a multiplicity of physiological roles, among them in the process of blood clot formation. It was found that it affects the metabolism of triacylglycerol-rich lipoproteins, the function of thrombocytes, and the activation of endothelial cells ([1] and references therein) and inhibits the transformation of prothrombin into thrombin [2]. It binds to structures which contain negatively charged phospholipid molecules such as phospholipid vesicles [3,4], thrombocytes [5], thrombocyte-derived microvesicles and apoptotic cells [6], serum lipoproteins [7,8] and mediates cellular recognition of negatively charged phospholipid-exposing microparticles [9-11].

II. METHODS

A. β 2GPI antibodies

β 2GPI (Hyphen BioMed, France) was aliquoted and stored at -70°C . In all experiments, the final concentration of β 2GPI in phosphate buffer saline (PBS) was 100 mg/L, which is approximately half the concentration of physiological β 2GPI in normal human plasma (about 200 mg/L) [12,7].

B. Giant phospholipid vesicles

GPVs were prepared at room temperature (23°C) by the electroformation method [13] modified as described in

Tomšič et al. [14]. The synthetic lipids cardiolipin (1,1'2,2'-tetraoleoyl cardiolipin), POPC (1-palmitoyl-2-oleoyl-sn-glycero-3-phosphocholine), and cholesterol were purchased from Avanti Polar Lipids, Inc. Appropriate volumes of POPC, cardiolipin and cholesterol, all dissolved in a 2:1 chloroform/methanol mixture, were combined in a glass jar and thoroughly mixed. For charged cardiolipin vesicles, POPC, cholesterol and cardiolipin were mixed in the proportion 2:2:1. For neutral POPC vesicles, POPC and cholesterol were mixed in proportion 4:1. Cholesterol was added to POPC to increase the longevity of vesicles. 10 μL of lipid mixture was applied to platinum electrodes. The solvent was allowed to evaporate in a low vacuum for 2 hours. The coated electrodes were placed in the electroformation chamber which was then filled with 3 mL of 0.2 M sucrose solution. An AC electric current with an amplitude of 5 V and a frequency of 10 Hz was applied to the electrodes for 2 hours, which was followed by 2.5 V and 5 Hz for 15 minutes, 2.5 V and 2.5 Hz for 15 minutes and finally 1 V and 1 Hz for 15 minutes. The content was rinsed out of the electroformation chamber with 5 mL of 0.2 M glucose and stored in a plastic test tube. The vesicles were left for sedimentation under gravity for one day at 4°C . 200 to 400 μL of the sediment was collected from the bottom of the tube and used for a series of experiments. Before placing the vesicles into the observation chamber, the sample was gently mixed.

C. Mathematical model

The mechanical behavior of the vesicle biomembrane is described using an elastic continuum theory where the bindings between the building blocks of the membrane (phospholipids) are represented as the elastic springs. For the sake of simplicity, two dimensional problem was considered. Energetically optimal deformed shape is computed via minimization of the elastic energy of a discrete linear model of lipids interaction.

The model of coalescence of phospholipid vesicles is based on the equilibrium between the two energy contributions: the adhesion energy (E_a) due to the contact, and the

bending and stretching elastic energy (E_e) of the membrane due to deformation. The adhesion energy was taken to be depended on the contact surface area linearly ($E_a = \gamma A$). As the deformation is prescribed by the shape of the contact area, the elastic energy is linearly dependent on an elastic constant, K . Therefore, the total energy of the system is computed as the energy difference between the elastic and the adhesion energy normalized to the elastic constant (K). Coefficient Γ describes dependence of the adhesion energy on the dimensionless contact area.

It was considered that the mother and bud vesicles have a form of sphere before adhesion with the radius of $1\ \mu\text{m}$ and $0.2\ \mu\text{m}$, respectively. The volume of the bud vesicle was taken as fixed. The thickness of the membrane was $5\ \text{nm}$.

III. RESULTS AND DISCUSSION

β 2GPI caused coalescence of cardiolipin-containing as well as of POPC vesicles. Adhesion to the bottom of the observation chamber occurred simultaneously. Formation of sticky complexes was also observed in the sample containing both kinds of vesicles. This indicates that β 2GPI mediates the interaction between the charged – charged, charged – neutral and neutral – neutral pairs of membranes.

If β 2GPI was present in the solution, the bud (Fig.1A-C) had coalesced with the mother vesicle before it could detach from it (Fig.1D-F).

Fig. 2 presents dependence of the ratio between the energy and the elastic constant on the magnitude of the dimensionless coalescence area for various coefficient of adhesion (Γ). Increasing the area of the contact, the elastic energy is increasing while the adhesion energy decreases. Equilibrium of the system is characterized by the minimum in the total energy. It is obvious from Fig. 2 that the area of contact in equilibrium state depends on the adhesion coefficient:

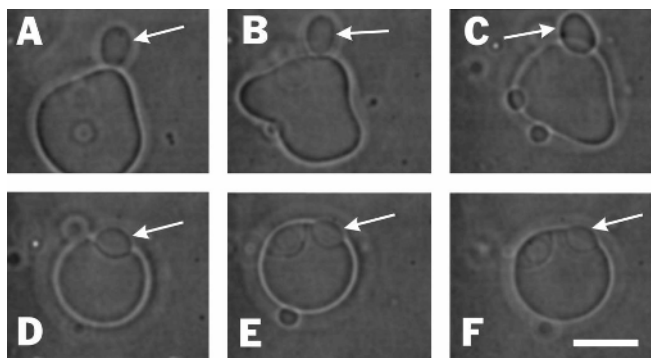


Fig. 1: The effect of β 2GPI dissolved in phosphate buffer saline (PBS) on a budding vesicle (A-F). The bud (marked by a white arrow) coalesced with the mother vesicle and remained attached to it. Bar denotes $10\ \mu\text{m}$.

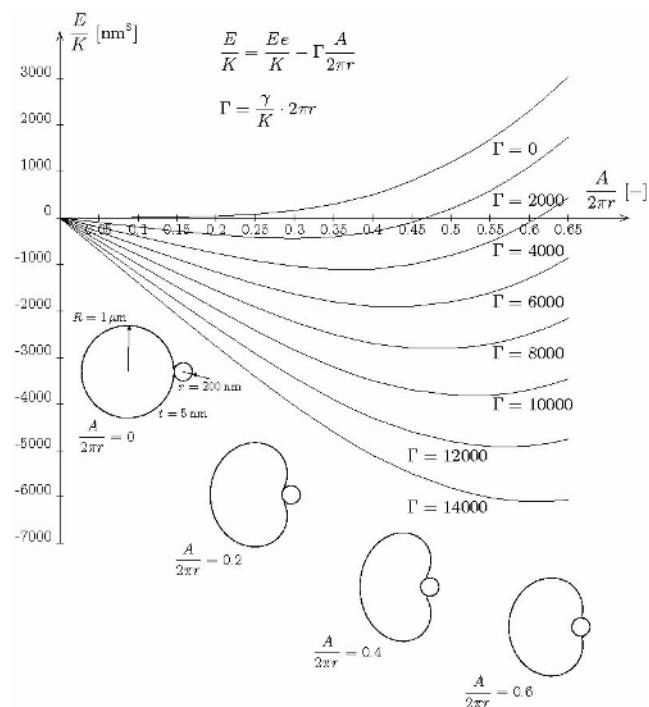


Fig. 2: Dependence between the dimensionless contact area (A) and total energy E normalized to the elastic constant (K). The equilibrium shapes of mother vesicle for various contact areas are shown.

coefficient: i.e. is zero for no adhesion ($\Gamma=0$) and large for strong adhesion ($\Gamma=14000\ \text{nm}^3$). The calculated shapes of equilibrium states correspond with the shapes observed in experiment.

IV. CONCLUSIONS

The presented model provides a simplified analysis of the problem neglecting the liquidity of the bilayer and reducing the problem to two dimensions. However, good agreement between the experimental measurements and the model simulations show that the phospholipid vesicle shapes during the coalescence of vesicles mediated by β 2GPI may be explained as an interplay between the deformation and contact energy. To describe this phenomenon in details, the mathematical model should be developed further.

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Author: Jasna Urbanija
 Institute: Laboratory of Clinical Biophysics,
 Faculty of Medicine, University of Ljubljana
 Street: Lipičeva 2
 City: Ljubljana
 Country: Slovenia
 Email: jasna.urbanija@fe.uni-lj.si